

#### Structure of the presentation

Part I

A brief description of the Institute for Problems in Mechanics (IPM) of Russian Academy of Sciences

#### Part II

An introduction to terminology and basic problems in the Surface Acoustic Wave (SAW) studies

#### Part III

Some recent results in SH- and Love wave analyses

#### Part I

# Institute for Problems in Mechanics of Russian Academy of Sciences

### History of the IPM



The first director of the former Institute of Mechanics Professor **B.G. Galerkin** ~1938



The first director of the Institute for Problems in Mechanics of Russian Academy of Sciences **A.Y. Ishlinskiy** 1965-1990

### History of the IPM

The following persons worked in IPM in different years

Professor L.A. Goldenveizer (theory of shells, asymptotic methods)

Professor N.V. Zvolinskij (theory of surface waves, analysis of tsunami)

Professor G.I. Barenblatt (fracture mechanics, hydrodynamics)

Professor O.A. Olejnik

(differential equations, homogenization)



### Laboratory of Wave Dynamics



#### The head of Laboratory Professor I.V. Simonov

#### Fracture mechanics and Astromechanics





### Laboratory of Wave Dynamics

The main directions:

Analyses of bodies collisions and aftereffects: Fragmentation, Cavern-and-crack formation, Penetration





#### Laboratory of Wave Dynamics

#### **Experimental study:**

Penetrating ability of different arrow-heads



The winner with respect to the penetrating ability



#### Part II

# Introduction to the Surface Acoustic Wave analyses

# The main types of Surface Acoustic Waves

### Rayleigh waves



### **Basic definitions**

Pioneering Rayleigh's work (1885) where these waves were described.

Rayleigh surface wave means attenuating with depth elastic wave with a plane wave front propagating on a traction-free boundary of a half-space.

# Rayleigh waves, their role in seismic activity

These waves play a very important role in transmitting the seismic energy and causing the catastrophic destructions due to the seismic activity.



# Rayleigh waves, danger for structures





### Rayleigh waves



# Problem of finding the secular equation

For an isotropic medium the secular equation (Rayleigh equation) is:

$$(\lambda + 2\mu)(\lambda + \mu)x^{3} - 8\mu(\lambda + 2\mu)(\lambda + \mu)x^{2} + 8\mu^{2}(\lambda + \mu)(3\lambda + 4\mu)x - 16\mu^{3}(\lambda + \mu)^{2} = 0$$

where

What is the secular equation for an anisotropic medium?

### Rayleigh waves



# Problem of finding the

#### secular equation

For an arbitrary anisotropic medium the secular equation has not been found yet (2006)

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Currie,	1979
Taziev,	1989
Destrade,	2001
Ting,	2002, 200

**Related problem** 







For a long time the main problem related to Rayleigh wave propagating on an **arbitrary anisotropic** half-space was finding conditions at which such a wave cannot propagate (Problem of "forbidden" directions).

**Do "forbidden directions" exist?** 



Rayeleigh waves

# The main practical problem

Theorem of existence for Rayleigh waves:

Barnett and Lothe,	1973-76
Chadwick,	1975-85
Ting,	1983-96

In 1998-2002 a type of Non-Rayleigh waves was observed and constructed explicitly.

Such a wave corresponds to appearing the Jordan blocks in a six-dimensional matrix associated with the Christoffel equation.

# A remark on Jordan blocks appearing in the matrix analysis



Marie Ennemond Camille Jordan 1838 - 1922



#### Lamb waves

# **Basic definitions**



These waves, discovered by Horace Lamb (1917), can propagate in a layer with either traction-free, clamped or mixed boundary conditions, imposed on the outer surfaces of a layer.



#### Lamb waves



### **Basic** properties

In contrast to Rayleigh waves, Lamb waves are highly dispersive, that means the the phase speed depends upon frequency or wavelength.

There can be an infinite number of Lamb waves propagating with the same phase speed and differing by the frequency.

Lamb waves can travel with both sub, intermediate, and supersonic speed.

An interesting physical observation:

After excitation, the most energy is transferred by the two lowest modes (symmetric and flexural).

#### Lamb waves

### The main problems



Presumably, the main problem related to the genuine Lamb wave propagation lies in constructing solutions for an anisotropic layer having arbitrary elastic anisotropy:

Deriving explicit secular equation(s) for the phase speed and frequency (for different speed intervals there may be needed different equations)

Obtaining and analyzing solutions for waves related to appearing the Jordan blocks in a sixdimensional matrix formalism.

### Stoneley waves



### Definition

These waves were described by Robert Stoneley (1924), and they are waves traveling on an interface between two contacting half-spaces.



### Stoneley waves



### **Basic** properties

Stoneley waves exponentially attenuate with depth in both half-spaces, and in this respect resemble Rayleigh waves.

Both Rayleigh and Stoneley waves are not dispersive.

#### Remark

For contacting isotropic half-spaces conditions of existence were found by Stoneley.

Conditions for existence of Stoneley waves propagating in anisotropic half-spaces are not established (2006).

### **Robert Stoneley**



### A fact from biography



#### Love waves



### Definition

These waves originate to Augustus Love (1911), who for the first time obtained a solution for a wave traveling in a system consisted of a layer and a contacting half-space.



#### SH waves

No image available

### Definition

These waves travel in a layer or possibly several contacting layers, and have the transverse horizontal polarization.





### Main definition

Love wave is a surface wave having horizontal transverse polarization and propagating in a medium composed of an elastic layer lying on a substrate.

#### MORE INFO...

It is assumed that Love wave attenuates with depth in a substrate.

## Love waves: polarization



### **Basic theoretical works**

Love	1911	<b>Representation</b> for surface waves with SH-
2010		polarization, propagating in a system
		composed of an <i>isotropic</i> layer lying on
		<i>isotropic</i> substrate
Thomson	1950	The first analysis of waves in stratified media
Haskell	1953	Correction of the previous results,
		introduction of the Transfer Matrix Method
Knopoff	1964	Introduction of the Global Matrix Method
Dieulesaint	1980	Equations for speed and polarization of
et Royer		Love waves in <i>orthotropic</i> media

### Multilayered structure



# The main problem of practical importance

Determining geometrical and physical properties of the internal layer(s) by analyzing the dispersion relations of Love surface waves propagating in the multi layered structure

### Analysis scheme

#### I. Single layer

- i. Obtaining Christoffel equation;
- ii. Determining the Christoffel parameters

#### II. Multiple layers

- i. Formulating contact type boundary conditions
- ii. Obtaining the Global (Transfer) Matrix
- **III.** Numerical implementation

Single layer analysis. Equations of motion

# $\mathbf{A}(\partial_x, \partial_t)\mathbf{u} \equiv \operatorname{div}_x \mathbf{C} \cdot \nabla_x \mathbf{u} - \rho \ddot{\mathbf{u}} = 0$

#### WHERE

- Is the elasticity tensor, assumed to be positive definite and hyper elastic;
- **U** Is the displacement field;
- $\rho$  Is the material density;

## Single layer analysis. Monoclinic anisotropy

#### **Definition:**

The material is called monoclinic (with respect to a direction **m**) if its symmetry group is generated by  $\mathbf{R}_{\mathbf{m}}^{\pi}$ 

#### REMARKS

- The definition is equivalent to vanishing all of the decomposable components of the tensor C having the odd number of entries of vector m;
- II. The assuming monoclinic symmetry provides a sufficient condition for the surface tractions acting on any plane (v ·x=const) to be collinear with vector m;
- III. The elasticity tensor for monoclinic medium has 13 independent elasticity constants

Single layer analysis. Representation for displacements for Love wave

 $\mathbf{u}(\mathbf{x},t) \equiv \mathbf{m} f(ir x') e^{ir(\mathbf{n} \cdot \mathbf{x} - ct)}$ 

#### WHERE

 $\mathbf{m} = \mathbf{v} \times \mathbf{n} \quad \text{is the amplitude vector ;} \\ \mathbf{n} \quad \text{is direction of propagation;} \\ \mathbf{v} \quad \text{is the unit normal to the median plane;} \\ x' = \mathbf{v} \cdot \mathbf{x} \\ f \quad \text{is the unknown scalar function;} \\ 3/13/2009$ 

Single layer analysis. The Christoffel equation  $\begin{pmatrix} (\mathbf{m} \otimes \mathbf{v} \cdot \mathbf{C} \cdot \mathbf{v} \otimes \mathbf{m}) \partial_{x'}^2 + (\mathbf{m} \otimes \mathbf{v} \cdot \mathbf{C} \cdot \mathbf{n} \otimes \mathbf{m} + \mathbf{m} \otimes \mathbf{n} \cdot \mathbf{C} \cdot \mathbf{v} \otimes \mathbf{m}) \partial_{x'} + \\ (\mathbf{m} \otimes \mathbf{n} \cdot \mathbf{C} \cdot \mathbf{n} \otimes \mathbf{m} - \rho c^2) \end{pmatrix} f(x') = 0$ The corresponding characteristic equation is:  $\left( (\mathbf{m} \otimes \mathbf{v} \cdots \mathbf{C} \cdots \mathbf{v} \otimes \mathbf{m}) \gamma^{2} + (\mathbf{m} \otimes \mathbf{v} \cdots \mathbf{C} \cdots \mathbf{n} \otimes \mathbf{m} + \mathbf{m} \otimes \mathbf{n} \cdots \mathbf{C} \cdots \mathbf{v} \otimes \mathbf{m}) \gamma + (\mathbf{m} \otimes \mathbf{n} \cdots \mathbf{C} \cdots \mathbf{n} \otimes \mathbf{m} - \rho c^{2}) \right) = 0$ 

#### REMARK

Multiple roots arise when the dicriminant vanishes

Single layer analysis. <u>Representations for the solution</u>

#### Aliquant roots

 $\mathbf{u}(\mathbf{x},t) = \mathbf{m} \left( C_1 \sinh(ir\alpha x') + C_2 \cosh(ir\alpha x') \right) e^{ir(\beta x' + \mathbf{n} \cdot \mathbf{x} - ct)}$ 

#### WHERE

 $C_1, C_2$  are arbitrary coefficients defined by boundary and interfacial conditions

# Boundary conditions Single layer on a substrate

I. Traction-free boundary conditions at the upper boundary h

$$\mathbf{t}_{layer}(\frac{h}{2},t) = 0$$

II. Contact type boundary conditions at the interface

$$\mathbf{u}_{layer}(-\frac{h}{2},t) = \mathbf{u}_{substrate}(0,t)$$
$$\mathbf{t}_{layer}(-\frac{h}{2},t) = \mathbf{t}_{substrate}(0,t)$$

III. Sommerield's attenuation condition

$$\lim_{\alpha' \to -\infty} \mathbf{u}_{substrate}(x',t) = 0$$



Some analytical results for Love waves (one orthotropic layer on the orthotropic substrate)

Proposition. a) No Love wave can propagate in a system composed of a single orthotropic layer lying on an orthotropic substrate, when multiple roots in the Christoffel equation for the layer arise;

b) Love wave can propagate, if and only if the phase speed belongs to the interval c∈(c<sub>layer</sub><sup>bulk</sup>; c<sub>substr</sub><sup>bulk</sup>);

c) The dispersion relation admits the following representation:

$$\omega = \frac{c}{\gamma_1 h_1} \left( \arctan \left( i \frac{\gamma_2 \left( \mathbf{m} \otimes \mathbf{v} \cdot \cdot \mathbf{C}_2 \cdot \cdot \mathbf{v} \otimes \mathbf{m} \right)}{\gamma_1 \left( \mathbf{m} \otimes \mathbf{v} \cdot \cdot \mathbf{C}_1 \cdot \cdot \mathbf{v} \otimes \mathbf{m} \right)} \right) + n\pi \right), \quad n = 0, 1, 2, ...;$$

### Love waves (one layer on a substrate)

20

16

Erequency 8

4

0.0



3/13/2009

A layer should be less rigid, than a substrate. Otherwise, Love waves do not exist.

1.0

**Dispersion curves** Love waves: one layer on a substrate (isotropic components,  $h_1=1$ ;  $\rho_1=\rho_2=1$ ,  $\mu_1=1$ ;  $\mu_2=4$ )

2.0

Phase speed

3.0

4.0



### (two layered plate)





# Love and SH waves (comparison)





#### Love and SH waves



### Love waves in multilayered medium



Love waves in multilayered medium. Dispersion curves (10-layered plate)



### Love waves in multilayered medium. Dispersion curves (10-layered plate)



### Love waves in multilayered medium. Layer thickness variation

Variation of lower branch of the dispersion curve due to 10% depth increase of the corresponding layer



Love waves in multilayered medium. Simultaneous variation of density and shear modulus of the corresponding layer/substrate



Love waves in multilayered medium. Delamination of the 10-layered plate from the substrate



Love waves in multilayered medium. Defoliation of the 10-layered plate from the substrate



# Numerical analysis: SH waves in stratified plates

SH waves in multilayered plates Lower mode dispersion curves for tractionfree plates with different number of layers



Plates with alternating isotropic layers:

$$h_1 = ... = h_n = 1$$
  
 $\rho_1 = ... = \rho_n = 1$   
 $\mu_1 = 1; \ \mu_2 = 4;$   
 $\mu_3 = 1; ...$ 

### SH waves in multilayered plates

A 31-layered plate

### SH waves in multilayered plates



### SH waves in multilayered plates



Possible applications of analyses for Love and SH waves propagating in stratified media

### **Possible** applications

Determination of physical properties of the internal layer(s) with questionable properties by analyzing the dispersion relations for Love/SH waves



### **Possible** applications

Application to geomechanics



Principle possibility to analyze depth, geometrical and physical properties of the questionable layers (water or oil saturated) by the dispersion curve analysis

### **Possible** applications

Application to glue laminated timber structure analysis



Principle ability to analyze physical properties, presence of cracks, flaws, and delaminations by the dispersion curve analysis